

Fast Multipole Method Algorithm in 2D

— Input

N : number of particles \rightarrow int npar

\rightarrow double $z[\text{Max_particle}][2]$
 $\{z_j = x_j + iy_j \mid x_j, y_j \in [0, 1]; i = \sqrt{-1}; j = 0, \dots, N\}$: particle positions

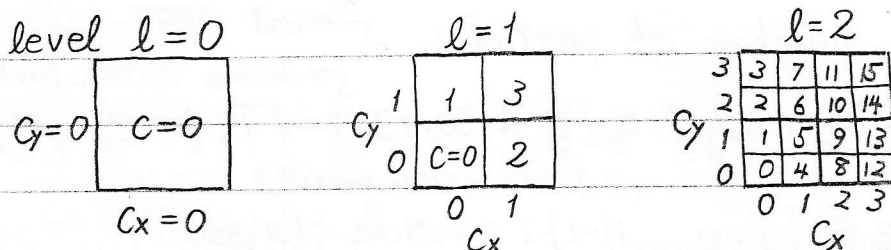
\rightarrow double $q[\text{Max_particle}]$
 $\{q_j \in \mathbb{R} \mid j = 0, \dots, N-1\}$: particle charges.

— Output

\rightarrow double pot[Max_npar][2]

$\{\phi(z_j) = \sum_{\substack{k=1 \\ (k \neq j)}}^N \frac{q_k}{z_j - z_k} \log(z_j - z_k) \mid j = 1, \dots, N\}$: electrostatic potentials

- Data structures: cell quadtree



\rightarrow int L

L: maximum level of refinement

Level $l=0$ is the root level, and the only one cell $C=0$ at $l=0$ is the entire simulation system $[0,1]^2$.

At level $l \geq 1$, each mother cell at level $l-1$ is subdivided into 2×2 square daughters of equal area, so that there are 4^l cells at level l . The recursive subdivision is repeated until the leaf level, $l=L$.

(Cell index)

At level l , the 4^l cells, $C=0, \dots, 4^l-1$, are indexed using a vector index, $\vec{C} = (C_x, C_y)$ ($C_x, C_y = 0, \dots, 2^l-1$), where C_x and C_y are the column and row indices in the x and y directions, respectively. The serial cell index is numbered in the column-major order:

$$\begin{cases} C_x = C / 2^l & \text{(quotient)} & (1) \\ C_y = C \bmod 2^l & \text{(remainder)} & (2) \end{cases}$$

or

$$C = C_x \cdot 2^l + C_y \quad (3)$$

- Data structures: multipoles and local expansions

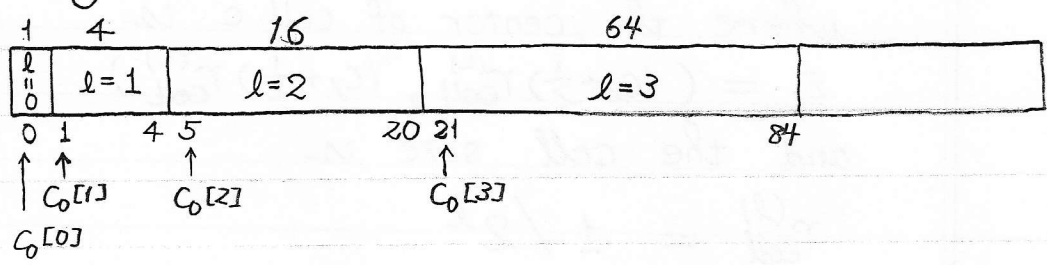
→ int P

P: Truncation degree of multipole and local expansions

$\{\Phi_c^{(l)}(\alpha) \mid l=0, \dots, L; c=0, \dots, 4^l-1; \alpha=0, \dots, P\}$: α -th order multipole around the center of cell c at level l
 ↳ phi[Max_cell][Max_term][Z]
 real(0) / imaginary(1)

$\{\Psi_c^{(l)}(\alpha) \mid l=0, \dots, L; c=0, \dots, 4^l-1; \alpha=0, \dots, P\}$: local expansion terms
 ↳ psi[Max_cell][Max_term][Z]

In the arrays phi & psi, cells are sequentially stored starting from 0th level.



$C_0(l)$ ($l=0, \dots, L$): the starting position, in arrays phi & psi, of the cell at level l.

↳ int CO[Max_level]

phi[$C_0(l)+c$][α][] stores $\Phi_c^{(l)}(\alpha)$
 psi[] stores $\Psi_c^{(l)}(\alpha)$

$$C_0(l) = \begin{cases} 0 & (l=0) \\ 1 + 4^k + \dots + 4^{l-1} = \frac{4^l - 1}{3} & (l \geq 1) \end{cases} \quad (4)$$

- Algorithm

1. Form the multipoles of all leaf cells at level L

for $c = 0$ to $4^L - 1$

$$\Phi_c^{(L)}(0:P) \leftarrow 0$$

for $j = 0$ to $N-1$ // scan \forall particles

$$c_x \leftarrow \lfloor x_j / r_{\text{cell}}^{(L)} \rfloor; c_y \leftarrow \lfloor y_j / r_{\text{cell}}^{(L)} \rfloor \quad // \text{particle} \rightarrow \text{cell mapping}$$

$$c \leftarrow c_x \cdot 2^L + c_y$$

$$\Phi_c^{(L)}(\alpha) += \begin{cases} q_j & (\alpha=0) \\ -\frac{q_j (z_j - z_c)^\alpha}{\alpha} & (\alpha \geq 1) \end{cases} \quad (5)$$

where the center of cell c is

$$z_c = \left((c_x + \frac{1}{2}) r_{\text{cell}}^{(L)}, (c_y + \frac{1}{2}) r_{\text{cell}}^{(L)} \right) \quad (6)$$

and the cell size is

$$r_{\text{cell}}^{(L)} = 1 / 2^L \quad (7)$$

2. Upward pass to compute the multipoles

for $l = L-1$ down to 0

for $c = 0$ to $4^l - 1$

$$\Phi_c^{(l)} \leftarrow \sum_{c' \in \text{daughters}(c)} T_{\mathbb{Z}_{c'} - \mathbb{Z}_c}^{M \leftarrow M} \Phi_{c'}^{(l+1)} \quad (\alpha = 0, \dots, P) \quad (8)$$

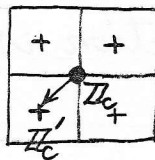
Here, $\text{daughters}(c)$ is the set of 4 daughter cells of cell c ,

$$\text{daughters}(c) = \left\{ c' \mid c'_\mu = 2c_\mu \text{ or } 2c_\mu + 1 \quad (\mu = x, y) \right\} \quad (9)$$

$$\mathbb{Z}_{c'} - \mathbb{Z}_c = \left(c'_x + \frac{1}{2}, c'_y + \frac{1}{2} \right) r_{\text{cell}}^{(l+1)} - \left(c_x + \frac{1}{2}, c_y + \frac{1}{2} \right) r_{\text{cell}}^{(l)} \quad (10)$$

and the multipole-to-multipole transformation is

$$T_{\mathbb{Z}}^{M \leftarrow M} \Phi_{c'}^{(l+1)}(\alpha) = \begin{cases} \Phi_{c'}^{(l+1)}(0) & (\alpha = 0) \\ \left[\sum_{\beta=1}^{\alpha} \Phi_{c'}^{(l+1)}(\beta) \bar{z}^{\alpha-\beta} C_{\beta-1} \right] - \frac{\Phi_{c'}^{(l+1)}(0) \bar{z}^{\alpha}}{\alpha} & (\alpha \geq 1) \end{cases} \quad (11)$$



3. Downward pass to compute the local expansion

for $c = 0$ to 4^l

$$\bar{\Psi}_c^{(2)}(0:P) \leftarrow 0$$

for $l = 2$ to L

for $c = 0$ to 4^l

$$\bar{\Psi}_c^{(l)} \leftarrow T_{\mathbb{Z}_{\text{mother}(c)} - \mathbb{Z}_c}^{L \leftarrow L} \bar{\Psi}_{\text{mother}(c)}^{(l-1)}$$

$$+ \sum_{c' \in \text{interactive}(c)} T_{\mathbb{Z}_{c'} - \mathbb{Z}_c}^{L \leftarrow M} \bar{\Phi}_{c'}^{(l)} \quad (12)$$

Here, the mother cell of cell c is

$$\text{mother}(c) = (c_x/2, c_y/2) \quad (13)$$

and the set of interactive cells is

$$\text{interactive}(c) = \{c' \mid c'_\mu = 2(\underbrace{c_\mu/2}_{\text{mother}} - 1), \dots, 2(c_\mu/2 + 1) + 1 \ (\mu=x,y)\}$$

$$\text{but any } c'_\mu \neq c_\mu \text{ or } c_\mu \pm 1 \ (\mu=x,y) \} \quad (14)$$

The multipole-to-local transformation is

$$\begin{aligned} & T_{\mathbb{Z}}^{L \leftarrow M} \bar{\Phi}_{c'}^{(l)}(\alpha) \\ &= \begin{cases} \sum_{\beta=1}^P \frac{\bar{\Phi}_{c'}^{(l)}(\beta)}{z^\beta} (-1)^\beta + \bar{\Phi}_{c'}^{(l)}(0) \log(-z) & (\alpha=0) \\ \left[\frac{1}{z^\alpha} \sum_{\beta=1}^P \frac{\bar{\Phi}_{c'}^{(l)}(\beta)}{z^\beta} C_{\alpha+\beta-1} (-1)^\beta \right] - \frac{\bar{\Phi}_{c'}^{(l)}(0)}{\alpha z^\alpha} & (\alpha \geq 1) \end{cases} \quad (15) \end{aligned}$$

The local-to-local transformation is

$$T_{\mathbb{Z}}^{L \leftarrow L} \Psi_{c'}^{(l-1)}(\alpha) = \sum_{\beta=\alpha}^P \Psi_{c'}^{(l-1)}(\beta) \beta C_{\alpha} (-z)^{\beta-\alpha} \quad (16)$$

(Detailed pseudocode for Eq. (12))

$$c'_x \leftarrow c_x/2; \quad c'_y \leftarrow c_y/2; \quad c' \leftarrow c'_x \cdot 2^{l-1} + c'_y \quad // \text{mother}$$

$$z \leftarrow (c'_x + \frac{1}{2}, c'_y + \frac{1}{2}) r_{\text{cell}}^{(l-1)} - (c_x + \frac{1}{2}, c_y + \frac{1}{2}) r_{\text{cell}}^{(l)} \quad // \mathbb{Z}^{\text{mother}} - \mathbb{Z}^{\text{daughter}}$$

$$\Psi_c^{(l)}(\alpha) \leftarrow \sum_{\beta=\alpha}^P \Psi_{c'}^{(l-1)}(\beta) \beta C_{\alpha} (-z)^{\beta-\alpha} \quad // \text{inherit mother's local expansion} \\ (\alpha=0, \dots, P)$$

$$\text{for } c'_x = 2(c_x/2-1) \text{ to } 2(c_x/2+1) + 1$$

$$\text{for } c'_y = 2(c_y/2-1) \text{ to } 2(c_y/2+1) + 1$$

if $|c'_x - c_x| \geq 2$ or $|c'_y - c_y| \geq 2$ then

$$z \leftarrow (c'_x - c_x, c'_y - c_y) r_{\text{cell}}^{(l)}$$

for $\alpha = 0$ to P

$$\Psi_c^{(l)}(\alpha) += \begin{cases} \sum_{\beta=1}^P \frac{\Phi_{c'}^{(l)}(\beta)}{z^{\beta}} (-1)^{\beta} + \Phi_{c'}^{(l)}(0) \log(-z) & (\alpha=0) \\ \left[\frac{1}{z^{\alpha}} \sum_{\beta=1}^P \frac{\Phi_{c'}^{(l)}(\beta)}{z^{\beta}} \alpha + \beta - 1 C_{\beta-1} (-1)^{\beta} \right] - \frac{\Phi_{c'}^{(l)}(0)}{\alpha z^{\alpha}} & (\alpha \geq 1) \end{cases}$$

// interactive-cell contribution

4. Direct calculation of nearest-neighbor leaf-cell contributionfor $j = 0$ to $N-1$

$$\begin{aligned} \Phi(z_j) \leftarrow & \sum_{\alpha=0}^P \Psi_{c(j)}^{(L)}(\alpha) (z_j - z_{c(j)})^\alpha \\ & + \sum_{\substack{j' \in nn(c(j)) \\ j' \neq j}} q_{j'} \log(z_j - z_{j'}) \end{aligned} \quad (17)$$

Here, $c(j)$ is the ^{leaf} cell that particle j belong to

$$c_x = \lfloor x_j / r_{cell}^{(L)} \rfloor, \quad c_y = \lfloor y_j / r_{cell}^{(L)} \rfloor$$

and $nn(c(j))$ is the 27 nearest-neighbor leaf cells of cell $c(j)$, including $c(j)$ itself.

The first term in Eq.(17) is the local expansion of the potential from non-nearest-neighbor leaf cells; the second term is the direct sums over particles in the nearest-neighbor leaf cells. Use the linked-list cell method in `pmd.c` to calculate the second contribution.

Slightly Modified Presentation of Downward Pass

(6)

3. Downward pass to compute the local expansion

for $C = 0$ to 4^l

$$\Psi_c^{(1)}(0:P) \leftarrow 0$$

for $l = 2$ to L

for $C = 0$ to 4^l

$$\Psi_c^{(l)} \leftarrow T_{\mathbb{Z}_{\text{mother}(c)} - \mathbb{Z}_c}^{L \leftarrow L} \Psi_{\text{mother}(c)}^{(l-1)} \quad (12)$$

for $C = 0$ to 4^l

$$\Psi_c^{(l)} \leftarrow \Psi_c^{(l)} + \sum_{C' \in \text{interactive}(c)} T_{\mathbb{Z}_{C'} - \mathbb{Z}_c}^{L \leftarrow M} \Phi_{C'}^{(l)} \quad (12')$$

Here, the mother cell of cell c is

$$\text{mother}(c) = (C_x/2, C_y/2) \quad (13)$$

and the set of interactive cells is

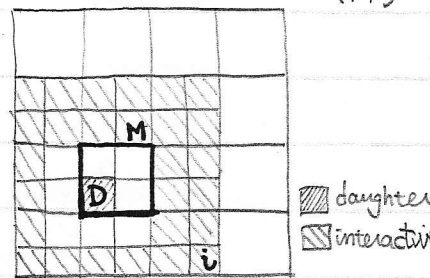
$$\text{interactive}(c) = \{ C' \mid C'_\mu = 2(\underbrace{C_\mu/2}_{\text{mother}} - 1), \dots, 2(C_\mu/2 + 1) + 1 \quad (\mu = x, y)$$

$$\text{but } |C'_x - C_x| > 1 \text{ or } |C'_y - C_y| > 1 \} \quad (14)$$

The multipole-to-local transformation is

$$T_{\mathbb{Z}}^{L \leftarrow M} \Phi_{C'}^{(l)}(\alpha)$$

$$= \begin{cases} \sum_{\beta=1}^P \Phi_{C'}^{(l)}(\beta) \left(-\frac{1}{\mathbb{Z}}\right)^\beta + \Phi_{C'}^{(l)}(0) \log(-\mathbb{Z}) & (\alpha=0) \\ \left[\frac{1}{\mathbb{Z}^\alpha} \sum_{\beta=1}^P \Phi_{C'}^{(l)}(\beta) \alpha + \beta - 1 C_{\beta-1} \left(-\frac{1}{\mathbb{Z}}\right)^\beta \right] - \frac{\Phi_{C'}^{(l)}(0)}{\alpha \mathbb{Z}^\alpha} & (\alpha \geq 1) \end{cases} \quad (15)$$



The local-to-local transformation is

$$T_{\mathbb{Z}}^{L \leftarrow L} \Psi_{C'}^{(l-1)}(\alpha) = \sum_{\gamma=0}^{P-\alpha} \Psi_{C'}^{(l-1)}(\alpha+\gamma) \alpha + \gamma C_\gamma (-\mathbb{Z})^\gamma \quad (16)$$

(Detailed pseudocode for Eq. (12))

$$C'_x \leftarrow C_x/2 ; C'_y \leftarrow C_y/2 ; C' \leftarrow C'_x \cdot 2^{l-1} + C'_y \quad // \text{mother}$$

$$\bar{z} \leftarrow (C'_x + \frac{1}{2}, C'_y + \frac{1}{2}) r_{\text{cell}}^{(l-1)} - (C_x + \frac{1}{2}, C_y + \frac{1}{2}) r_{\text{cell}}^{(l)} \quad // \mathbb{Z}^{\text{mother}} - \mathbb{Z}^{\text{daughter}}$$

$$\Psi_c^{(l)}(\alpha) \leftarrow \sum_{\gamma=0}^{P-\alpha} \Psi_{c'}^{(l-1)}(\alpha+\gamma) C_\alpha (-\bar{z})^\gamma \quad // \text{inherit mother's local expansion} \\ (\alpha=0, \dots, P)$$

(Detailed pseudocode for Eq. (12'))

$$\text{for } C'_x = C'_{x\text{-begin}} \quad \text{to} \quad C'_{x\text{-end}}$$

$$\text{for } C'_y = C'_{y\text{-begin}} \quad \text{to} \quad C'_{y\text{-end}}$$

if $|C'_x - C_x| > 1$ or $|C'_y - C_y| > 1$ then

} // interactive cells

$$\bar{z} \leftarrow (C'_x - C_x, C'_y - C_y) r_{\text{cell}}^{(l)}$$

for $\alpha = 0$ to P

$$\Psi_c^{(l)}(\alpha) += \begin{cases} \sum_{\beta=1}^P \Phi_{c'}^{(l)}(\beta) \left(-\frac{1}{\bar{z}}\right)^\beta + \Phi_{c'}^{(l)}(0) \log(-\bar{z}) & (\alpha=0) \\ \left[\frac{1}{\bar{z}^\alpha} \sum_{\beta=1}^P \Phi_{c'}^{(l)}(\beta) \alpha + \beta - 1 C_{\beta-1} \left(-\frac{1}{\bar{z}}\right)^\beta \right] - \frac{\Phi_{c'}^{(l)}(0)}{\alpha \bar{z}^\alpha} & (\alpha \geq 1) \end{cases}$$

where

$$C'_{\mu\text{-begin}} = \max(2(C_\mu/2 - 1), 0)$$

$$C'_{\mu\text{-end}} = \min(2(C_\mu/2 + 1) + 1, 2^l - 1) \quad (\mu = x, y)$$