

# Density Matrix Minimization: Real-Space Basis

6/19/03

[E. Hernandez, M.J. Gillan, C.M. Goringe, PRB 53, 7147 (1996)]

— Energy functional

$$\Omega = - \sum_i f_i \int dr \psi_i^*(r) (\hat{H}_r - \mu) \psi_i(r) \quad (1)$$

$$= \int dr [(\hat{H}_r - \mu) \sum_i f_i \psi_i(r) \psi_i^*(r')]_{r \rightarrow r}$$

$$\therefore \Omega = \int dr [(\hat{H}_r - \mu) \rho(r, r')]_{r \rightarrow r} \quad (2)$$

where

$$\rho(r, r') = \sum_i f_i \psi_i(r) \psi_i^*(r') \quad (3)$$

$f_i$  is the Fermi distribution,  $\{\psi_i(r)\}$  is the energy eigenstate set.

We generalize Eq. (2) to a functional of general density matrix  $\rho(r, r')$  with the following constraints:

(i) Normalization

$$N_e = \int dr \rho(r, r) \quad (4)$$

where  $N_e$  is the number of electrons.

(ii) Idempotency (at  $T=0$ )

$$\int dr' \rho(r, r') \rho(r', r'') = \rho(r, r'') \quad (5)$$

$$\hat{\rho}^2 = \hat{\rho} \quad (\text{see 6/18/03})$$

$$\langle r | \hat{\rho} | r' \rangle \langle r' | \hat{\rho} | r'' \rangle = \langle r | \hat{\rho} | r'' \rangle$$

$$\int dr' |r'\rangle \langle r'| = I \quad (\text{closure relation})$$

## ○ - Separable trial function

We restrict the variational space to the following separable form:

$$\rho(r; r') = \sum_{\alpha\beta} \phi_{\alpha}(r) K_{\alpha\beta} \phi_{\beta}^*(r') \quad (6)$$

↗ If this is diagonalized, SVD low-rank  
 ↘ ab initio basis to be determined (KS-like)

where  $\phi_{\alpha}(r)$  is a localized basis (support function). The index  $\alpha$  collectively denotes the atom index  $i$  and basis orbitals  $\mu$  per atom.  $\phi_{\alpha}$  are nonzero only inside spherical regions of radius  $R_{\text{reg}}$ .

$$\phi_{\alpha}(r) = 0 \quad (|r - R_{\alpha}| > R_{\text{reg}}) \quad (6)$$

where  $R_{\alpha}$  is the position of atom  $\alpha$ . (In DMDC, block size  $L_b$  becomes the corresponding parameter.)

## Density matrix purification

Given a trial density matrix

$$\sigma(r, r') = \sum_{\alpha\beta} \phi_{\alpha}(r) L_{\alpha\beta} \phi_{\beta}^*(r') \quad (7)$$

the purified density matrix, Eq. (6), is obtained as

$$\rho = 3\sigma * \sigma - 2\sigma * \sigma * \sigma \quad (8)$$

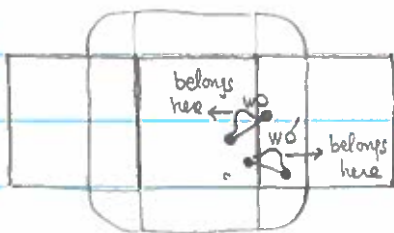
where  $C = A * B$  is a short-hand notation for

$$C(r, r') = \int dr'' A(r, r'') B(r'', r') \quad (9)$$

\* Note that purification only applies to subspace rotation,  $L \rightarrow K$ , but not to the ab initio orbitals  $\{\phi_{\alpha}(r)\}$ , which is the same in the LCWO in DMDC.

↓ Redundancy-ilconditioning avoidance

DMDC recipe: Each block only keeps Wannier orbitals whose centers-of-mass  $\in$  region.



Substituting the separable trial form, Eq. (7), in (8),

$$\begin{aligned}
 \rho(1,2) &= 3 \int d^3z \sum_{\alpha\beta} \phi_\alpha(1) L_{\alpha\beta} \phi_\beta^*(3) \sum_{\gamma\delta} \phi_\gamma(3) L_{\gamma\delta} \phi_\delta^*(2) \\
 &\quad - 2 \int d^3z \int d^3z' \sum_{\alpha\beta} \phi_\alpha(1) L_{\alpha\beta} \phi_\beta^*(3) \sum_{\gamma\delta} \phi_\gamma(3) L_{\gamma\delta} \phi_\delta^*(4) \sum_{\eta\nu} \phi_\eta(4) L_{\eta\nu} \phi_\nu^*(2) \\
 &= 3 \sum_{\alpha\beta} \sum_{\gamma\delta} \phi_\alpha(1) L_{\alpha\beta} S_{\beta\gamma} L_{\gamma\delta} \phi_\delta^*(2) \\
 &\quad - 2 \sum_{\alpha\beta} \sum_{\gamma\delta} \sum_{\eta\nu} \phi_\alpha(1) L_{\alpha\beta} S_{\beta\gamma} L_{\gamma\delta} S_{\delta\eta} L_{\eta\nu} \phi_\nu^*(2) \\
 &= 3 \sum_{\alpha\delta} \phi_\alpha(1) (LSL)_{\alpha\delta} \phi_\delta^*(2) \\
 &\quad - 2 \sum_{\alpha\nu} \phi_\alpha(1) (LSLSL)_{\alpha\nu} \phi_\nu^*(2)
 \end{aligned}$$

$$\therefore \rho(r,r') = \sum_{\alpha\beta} \phi_\alpha(r) (3LSL - 2LSLSL)_{\alpha\beta} \phi_\beta^*(r') \quad (10)$$

where

$$S_{\alpha\beta} = \int dr \phi_\alpha^*(r) \phi_\beta(r) \quad (11)$$

Comparison of Eqs. (6) and (10) yields

$$K = 3LSL - 2LSLSL \quad (12)$$

(Local approximation)

$$L_{\alpha\beta} = 0 \quad (|R_\alpha - R_\beta| > R_L) \quad (13)$$

where  $R_L$  is the off-diagonal density matrix cut-off.

## — Density-matrix minimization

The grand potential, Eq. (2), is minimized with respect to  $L_{\alpha\beta}$  and  $\Phi_{\alpha}(r)$ . For DMDC, only variation with respect to  $L_{\alpha\beta}$  is relevant, since the ab initio Wannier orbitals  $\Phi_{\alpha}(r)$  are fixed in the first (divide-and-conquer) phase and their mesh-level variation is irrelevant (too fine grained) for the construction of global density matrix.

(Variation with respect to  $K_{\alpha\beta}$ )

We may take a two-phase approach:

Phase 1: McWeely purification (fixed-point iteration) of DM

for step = 1 to  $n_{fp}$

$$K \leftarrow 3KSK - 2KSKSK$$

endfor.

Phase 2: Energy minimization (greedy <sup>conjugate gradient</sup> algorithm)

for step = 1 to  $n_{cg}$

$$K \leftarrow K - \tau \frac{\partial \Omega}{\partial K}$$

endfor



(6)

Substitute the trial function, Eq.(6), in Eq. (2)

$$\begin{aligned}\Omega[K] &= \int dr \sum_{\alpha\beta} [(\hat{H}_r - \mu) \phi_\alpha(r) K_{\alpha\beta} \phi_\beta^*(r')]_{r' \rightarrow r} \\ &= \sum_{\alpha\beta} K_{\alpha\beta} \underbrace{\int dr \phi_\beta^*(r) (\hat{H}_r - \mu) \phi_\alpha(r)}_{H'_{\beta\alpha}}\end{aligned}$$

$$\therefore \Omega[K] = \sum_{\alpha\beta} K_{\alpha\beta} H_{\beta\alpha} \quad (7)$$

where

$$H'_{\alpha\beta} = \int dr \phi_\beta^*(r) (\hat{H}_r - \mu) \phi_\alpha(r) \quad (8)$$

Taking the variation w.r.t.  $K$ ,

$$\delta\Omega = \sum_{\alpha\beta} \delta K_{\alpha\beta} H'_{\beta\alpha} \equiv \sum_{\alpha\beta} \delta K_{\alpha\beta} \frac{\partial \Omega}{\partial K_{\alpha\beta}}$$

$$\therefore \frac{\partial \Omega}{\partial K_{\alpha\beta}} = H'_{\beta\alpha} \quad (9)$$