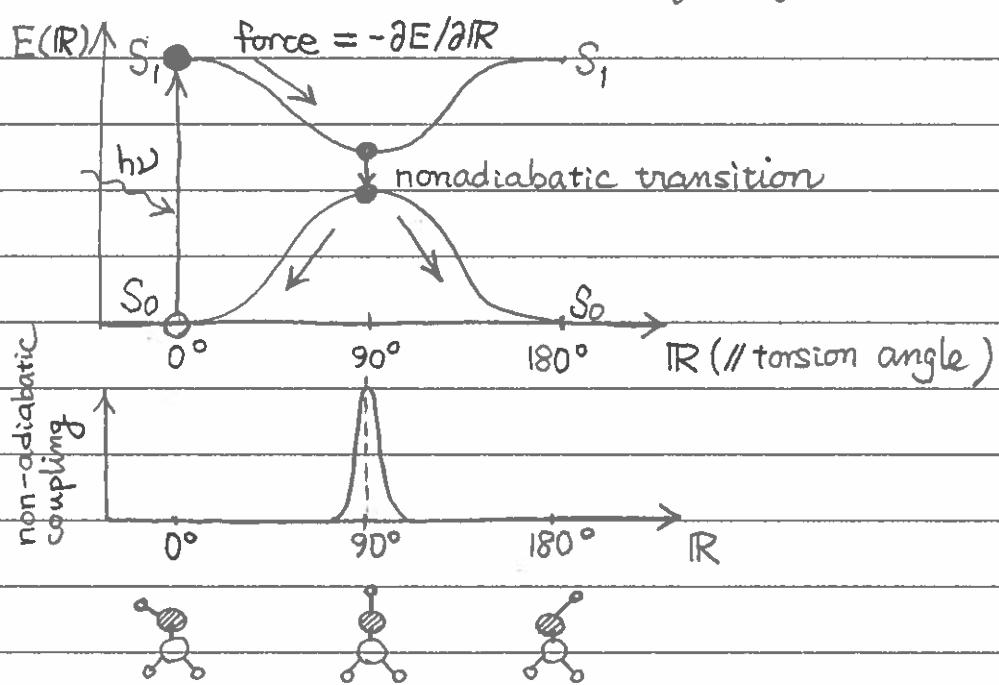


Excited-State Force Calculation Recipes

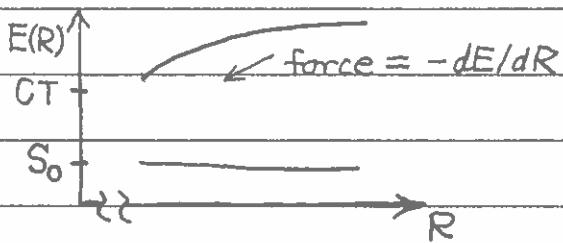
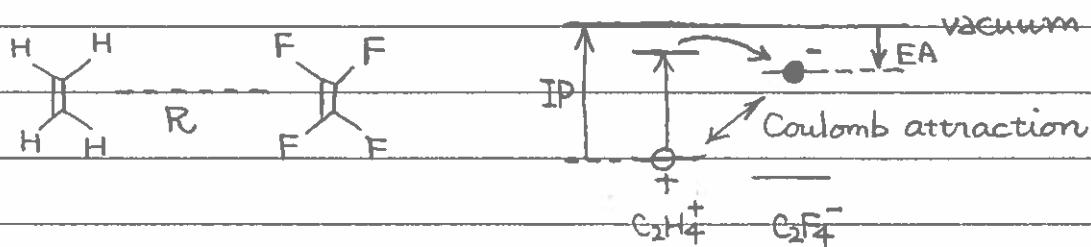
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- Goal: Low-cost, non-self-consistent (perturbative) calculation of forces on excited adiabatic surfaces, which captures correct physics.

(Example: Photo-isomerization of $\text{CH}_2=\text{NH}$)



(Example 2: Charge-transfer (CT) excitation)



(2)

The Coulomb attraction between donor⁺ and acceptor⁻ requires

- ① Self-consistent charge density / Kohn-Sham (KS) state calculation
w/ CT-excitation configuration.

OR

- ② Many-body (Casida) approach w/ ground-state KS basis, which eliminates self-interaction (so that the newly occupied acceptor state's energy E_{ao} does not include the Hartree-potential contribution from the evacuated $|\psi_{\text{lo}}(\mathbf{r})|^2$ in the donor).

- Finite-difference approach

1. Obtain the I-th excited-state energy, $E_I(R)$, as a linear combination of ground-state KS orbitals, using asymptotically correct (long-range exact exchange) exchange-correlation functional in Casida's equation.

2. Then compute the force by finite difference,

$$F = - \frac{E(R + \Delta) - E(R)}{\Delta}$$



Need 3N excited-energy calculations, which is prohibitive for large systems.

(3)

- Non-self-consistent (Harris-Foulkner) force calculation

1. Obtain ground-state KS orbitals self-consistently.

2. Solve Casida equation to obtain the I-th excited state,

$$\Psi_I = \sum_{ia\sigma} \frac{E_{ao} - E_{i\sigma}}{\omega_I} F_{ia\sigma}^I \hat{c}_{ao}^\dagger \hat{c}_{i\sigma} \Phi_0 \quad (1)$$

↓ Casida eigenvector ↑ Casida eigenvalue ↓ ground-state
 ↓ Slater determinant

OK to equate this = 1 for

large band-gap materials, if

keeping orthogonality is

beneficial [Walter & Häkkinen, New J. Phys. 10, 043018 ('08)]

3. Compute the density perturbation

$$\delta\rho(\mathbf{r}) = \rho[\mathbf{r}; \Psi_I] - \rho[\mathbf{r}; \Phi_0] \quad (2)$$

which contains the crucial information (e.g. electron-hole polarization for a CT excitation state).

4. Compute the excited-state force using the Harris-Foulkner perturbation approach [Torralba et al., JCTC 5, 1499 ('09)].

$$\mathbf{F}_n = \mathbf{F}_n^{\text{Hellmann-Feynman}} + \mathbf{F}_n^{\text{Non-selfconsistent}} \quad (3)$$

↓ force of n-th atom

where

(4)

$$\mathbb{F}_n^{\text{HF}} = - \int d\mathbf{r} \rho(\mathbf{r}) \frac{\partial V_{\text{loc}}}{\partial (\mathbf{r} - \mathbf{R}_n)} \quad (\text{local pseudo-potential})$$

↑ for Φ_I

$$= \sum_{ia\sigma} |\psi_{ia\sigma}^I|^2 \sum_{j\sigma} f_{ia\sigma}(\epsilon_j) \sum_{lm} \int d\mathbf{r} f_{lm} \left[\psi_{j\sigma}^*(\mathbf{r}) \frac{\partial \zeta_{lm}^i}{\partial (\mathbf{r} - \mathbf{R}_n)} \zeta_{lm}^i(\mathbf{r}' - \mathbf{R}_n) \psi_{j\sigma}(\mathbf{r}') \right]$$

weight of
each of
multi Slater-determinant
excited state

$$+ \psi_{j\sigma}^*(\mathbf{r}) \zeta_{lm}^i(\mathbf{r}' - \mathbf{R}_n) \frac{\partial \zeta_{lm}^i}{\partial (\mathbf{r}' - \mathbf{R}_n)} \psi_{j\sigma}(\mathbf{r}')$$

(nonlocal pseudo-potential)

(4)

$$\mathbb{F}_n^{\text{NSC}} = - \int d\mathbf{r} \left\{ S \mathcal{V}_{\text{Hartree}}(\mathbf{r}) \frac{\partial P[\mathbf{r}; \Phi_0]}{\partial \mathbf{R}_n} + S P(\mathbf{r}) \frac{\partial V_{xc}}{\partial \mathbf{R}_n} \right\} \quad (5)$$

(Please see Torralba et al. for the implementation of $\mathbb{F}_n^{\text{NSC}}$.)

- Test the NSC force calculation for the CT excitation state of $\text{C}_2\text{H}_4 - \text{C}_2\text{F}_4$ pair, and compare it with $-dE/dR$.