

Liouville Equation

[R. Zwanzig, "Nonequilibrium Statistical Mechanics", P.31]

- Phase space

Let $Q = (q_1, \dots, q_{3N})$ and $P = (p_1, \dots, p_{3N})$ be the coordinate and momentum vectors, where N is the number of atoms.

The pair of the vectors, $X = (Q, P)$, specifies a phase-space point.

- Phase-space trajectory: Liouville operator

The motion of the system in the phase space is described by the trajectory, $X(t)$ ($t \in \mathbb{R}$), which is governed by Hamilton's equations,

$$\begin{cases} \dot{q} = \frac{dQ}{dt} = \frac{\partial H}{\partial P} \end{cases} \quad (1a)$$

$$\begin{cases} \dot{p} = \frac{dP}{dt} = -\frac{\partial H}{\partial Q} \end{cases} \quad (1b)$$

For the Hamiltonian, $H(Q, P)$, we assume a form

$$H(Q, P) = \sum_{i=1}^{3N} \frac{P_i^2}{2m_i} + V(Q) \quad (2a)$$

$$= \frac{1}{2} P^T M^{-1} P + V(Q) \quad (2b)$$

where m_i is the mass associated with the i -th coordinate,

$M_{ij} = m_i \delta_{ij}$ is the diagonal mass matrix, and $V(Q)$ is the potential energy.

Substituting Eq. (2) in (1), the equations of motion become

$$\begin{cases} \dot{Q} = M^{-1} P \end{cases} \quad (3a)$$

$$\begin{cases} \dot{P} = -\frac{\partial V}{\partial Q} \end{cases} \quad (3b)$$

Consider an arbitrary function, $A(Q, P)$, of a phase-space point. The time evolution of the function value along the trajectory, X_t , is given by

$$\begin{aligned} \frac{dA}{dt} &= \dot{q} \cdot \frac{\partial A}{\partial q} + \dot{p} \cdot \frac{\partial A}{\partial p} \\ &= \underbrace{\frac{\partial H}{\partial p}}_{\dot{q}} \cdot \frac{\partial A}{\partial q} - \underbrace{\frac{\partial H}{\partial q}}_{\dot{p}} \cdot \frac{\partial A}{\partial p} \\ &= \left(\frac{\partial H}{\partial p} \frac{\partial}{\partial q} - \frac{\partial H}{\partial q} \frac{\partial}{\partial p} \right) A \end{aligned}$$

$$\therefore \frac{dA}{dt} = LA \quad (4)$$

where the Liouville operator is defined as

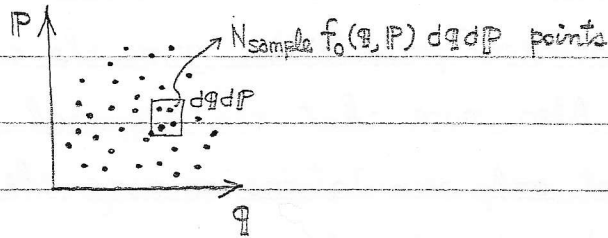
$$L = \frac{\partial H}{\partial p} \frac{\partial}{\partial q} - \frac{\partial H}{\partial q} \frac{\partial}{\partial p} \quad (5)$$

The formal solution of Eq.(4) is

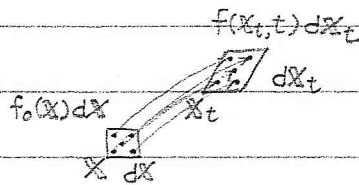
$$A(t) = e^{tL} A(p, q) \quad (6)$$

- Phase space distribution: Liouville's theorem.

Consider an ensemble of phase-space points, and let $f_0(Q, P) dQ dP = f_0(X) dX$ be the probability to find the system in a volume element dX around X .



Let all the phase-space points in the ensemble evolve in time, $X_t = e^{tL} X$, and consider the distribution, $f(X, t)$, at time t .



We consider a volume element, dX , around X , and follow the trajectories of $N_{\text{sample}} f_0(X) dX$ for time t , at which the points occupy a volume element, dX_t , around X_t :

$$\cancel{N_{\text{sample}} f(X_t, t) dX_t} = \cancel{N_{\text{sample}} f_0(X) dX} \quad (7)$$

$$f(X_t, t) \left| \frac{\partial X_t}{\partial X} \right| dX = f_0(X) dX \quad (8)$$

Now consider the Jacobian, $|\partial X_t / \partial X|$, in Eq. (8). For small time, dt , from Eq. (3),

$$\begin{cases} Q_{dt} = Q + \dot{Q} dt = Q + M^{-1} P dt & (9a) \end{cases}$$

$$\begin{cases} P_{dt} = P + \dot{P} dt = P - \frac{\partial V}{\partial Q} dt & (9b) \end{cases}$$

$$\begin{aligned}
 & \dots \begin{vmatrix} \partial \mathbb{P}_t / \partial \mathbb{Q} & \partial \mathbb{P}_t / \partial \mathbb{P} \\ \partial \mathbb{P}_t / \partial \mathbb{Q} & \partial \mathbb{P}_t / \partial \mathbb{P} \end{vmatrix} \\
 & = \begin{vmatrix} \mathbb{I} & M^{-1} dt \\ -\frac{\partial^2 V}{\partial \mathbb{Q}^2} dt & \mathbb{I} \end{vmatrix}
 \end{aligned}$$

Note that the identical permutation gives the contribution 1 to the determinant. Also note that the diagonal blocks are identity matrices, so that only permutations mixing the \mathbb{Q} and \mathbb{P} blocks need to be considered. Consider

$$J_{11} \ J_{22} \ \dots \ \underbrace{J_{\mathbb{Q}\mathbb{P}}}_{m_j^{-1} \delta_j} \ \dots \ \underbrace{J_{\mathbb{P}\mathbb{Q}}}_{\frac{\partial^2 V}{\partial \mathbb{Q}_j \partial \mathbb{Q}_j}} \ \dots \ J_{NN} = -\frac{1}{m_j} \frac{\partial^2 V}{\partial \mathbb{Q}_j \partial \mathbb{Q}_j} \delta_{ij} dt^2$$

Any such permutation thus introduces a factor dt^2 .

Therefore,

$$\left| \frac{\partial \mathbb{X}_{dt}}{\partial \mathbb{X}} \right| = 1 + \mathcal{O}(dt^2) \tag{10}$$

and

$$\begin{aligned}
 \left| \frac{\partial \mathbb{X}_t}{\partial \mathbb{X}} \right| &= \lim_{N \rightarrow \infty} \left(1 + a \left(\frac{t}{N} \right)^2 \right)^N \\
 &= \exp \left(\lim_{N \rightarrow \infty} N \log_e \left[1 + a \left(\frac{t}{N} \right)^2 \right] \right) \\
 &= \exp \left(\lim_{N \rightarrow \infty} N \cdot a \frac{t^2}{N^2} \right) = e^0 = 1
 \end{aligned}$$

$$\therefore \left| \frac{\partial \mathbb{X}_t}{\partial \mathbb{X}} \right| = 1 \tag{11}$$

Combining Eqs. (8) and (11),

$$f(\mathbb{X}_t, t) = f_0(\mathbb{X}) \tag{12}$$

(5)

For small dt , Eq. (12) results in

$$f(x + \dot{x} dt, dt) = f(x, 0)$$

$$f(x, dt) + \dot{x} dt \cdot \frac{\partial f}{\partial x} \Big|_{x, dt} = f(x, 0)$$

\dot{x} in the leading term

$$\begin{aligned} \frac{f(x, dt) - f(x, 0)}{dt} &= -\dot{x} \frac{\partial f}{\partial x} \\ &= -\left(\dot{q} \cdot \frac{\partial}{\partial q} + \dot{p} \cdot \frac{\partial}{\partial p} \right) f(x, 0) \\ &= -\left(\frac{\partial H}{\partial p} \frac{\partial}{\partial q} - \frac{\partial H}{\partial q} \frac{\partial}{\partial p} \right) f(x, 0) \end{aligned}$$

$$= -L f(x, 0) \tag{13}$$

$$\therefore \frac{\partial f}{\partial t} = \ominus L f \tag{14}$$

→ note the negative sign.

and its formal solution is

$$f(x, t) = e^{-tL} f(x, 0) \tag{15}$$

- Continuity equation

From Eq. (14)

$$\begin{aligned} \frac{\partial f}{\partial t} &= -\frac{\partial H}{\partial P} \cdot \frac{\partial f}{\partial Q} + \frac{\partial H}{\partial Q} \cdot \frac{\partial f}{\partial P} \\ &= -\frac{\partial}{\partial Q} \cdot \left(\frac{\partial H}{\partial P} f \right) + \frac{\partial}{\partial P} \cdot \left(\frac{\partial H}{\partial Q} f \right) - \frac{\partial^2 H}{\partial P \partial Q} f \\ &= -\frac{\partial}{\partial Q} \cdot (\dot{Q} f) - \frac{\partial}{\partial P} \cdot (\dot{P} f) \\ &= -\frac{\partial}{\partial X} \cdot (\dot{X} f) \end{aligned} \tag{16}$$

$$\therefore \frac{\partial f}{\partial t} + \frac{\partial}{\partial X} \cdot (\dot{X} f) = 0 \tag{17}$$

(Useful identity)

For arbitrary function, A(X),

$$LA(X) = \dot{X} \cdot \frac{\partial}{\partial X} A = \frac{\partial}{\partial X} \cdot (\dot{X} A) \tag{18}$$

$$\begin{aligned} \text{☺ } LA(X) &= \left(\frac{\partial H}{\partial P} \frac{\partial}{\partial Q} - \frac{\partial H}{\partial Q} \frac{\partial}{\partial P} \right) A = \dot{X} \cdot \frac{\partial}{\partial X} A \\ &= \frac{\partial}{\partial Q} \cdot \left(\frac{\partial H}{\partial P} A \right) - \frac{\partial}{\partial P} \cdot \left(\frac{\partial H}{\partial Q} A \right) + \frac{\partial^2 H}{\partial P \partial Q} A \\ &= \frac{\partial}{\partial X} \cdot (\dot{X} A) \end{aligned}$$