# **Unitary Time-Propagators**

Both space-splitting method (SSM) and spectral method (SM) are split-operator methods based on Trotter expansion. They perform *exact exponentiation* of an approximate time-propagation operator, by keeping infinite sum of Taylor expansion.

**Q.** Why exact exponentiation?

**A.** To conserve the probability.

# **Time-Propagation Operator**

 $|\psi(t)\rangle = e^{-iHt}|\psi(0)\rangle$ where  $|\psi(t)\rangle$  is wave function at time *t* and *H* is Hamiltonian operator.

#### **Norm Conservation**

$$\langle \psi(t)|\psi(t)\rangle = \langle \psi(0)|\underbrace{\left(e^{-iHt}\right)^{\dagger}}_{e^{iH^{\dagger}t}=e^{iHt}}e^{-iHt}|\psi(0)\rangle = \langle \psi(0)|\psi(0)\rangle = 1$$

Here, we have used Hermiticity of the Hamiltonian, *i.e.*, its Hermitian conjugate is itself.

### **Euler Method**

$$e^{-iH\Delta} = 1 - iH\Delta + O(\Delta^2)$$

 $\langle \psi(\Delta) | \psi(\Delta) \rangle = \langle \psi(0) | (1 + iH\Delta)(1 - iH\Delta) | \psi(0) \rangle = 1 + \Delta^2 \langle \psi(0) | H^2 | \psi(0) \rangle > 1$ Bad to truncate the Taylor expansion — probability not conserved!

### **Crank-Nicholson Method (CNM)**

$$e^{-iH\Delta} = \frac{1-iH\Delta/2}{1+iH\Delta/2} + O(\Delta^2)$$
  
$$\langle \psi(\Delta) | \psi(\Delta) \rangle = \langle \psi(0) | \frac{1+iH\Delta/2}{1-iH\Delta/2} \frac{1-iH\Delta/2}{1+iH\Delta/2} | \psi(0) \rangle = \langle \psi(0) | \psi(0) \rangle = 1 // \text{ It's unitary!}$$
  
In addition to SSM and SM, CNM is a unitary time propagator.

# **Application of rational operator**

$$|\psi(t)\rangle = \frac{1-iH\Delta/2}{1+iH\Delta/2}|\psi(0)\rangle \implies \underbrace{(1+iH\Delta/2)}_{A:\text{tridigonal matrix }x:\text{unknown vector}} \underbrace{|\psi(t)\rangle}_{b:\text{known vector}} = \underbrace{(1-iH\Delta/2)|\psi(0)\rangle}_{b:\text{known vector}}$$

Linear system of equations: Ax = b