

# Density Matrix Renormalization Group Revisited (1)

## Central Theorem - Optimal Low-Rank Density Matrix for a Block in Environment

6/16/03

### Definitions

Superblock ( $\mathcal{S}$ ) = a coupled block ( $\mathcal{B}$ ) + environment ( $\mathcal{E}$ ):

$$\mathcal{S} = \mathcal{B} \cup \mathcal{E} \quad (1)$$

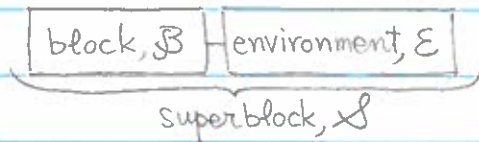
$$\mathcal{B} = \text{span}\{|i\rangle | i=1, \dots, N\} \quad (2)$$

$$\mathcal{E} = \text{span}\{|j\rangle | j=1, \dots, M\} \quad (3)$$

where  $\{|i\rangle\}$  and  $\{|j\rangle\}$  are orthonormal bases,

$$\langle i | i' \rangle = \delta_{ii'} \quad (4)$$

$$\langle j | j' \rangle = \delta_{jj'} \quad (5)$$



(Ground state)

$$|\Psi\rangle = \sum_{i=1}^N \sum_{j=1}^M \psi_{ij} |i\rangle |j\rangle \quad (6)$$

(Reduced density matrix)

Consider an arbitrary operator  $A$  that operates only in  $\mathcal{B}$ . Its expectation value is obtained as

$$\langle A(\mathcal{B}) \rangle = \sum_i \sum_j \langle j | \langle i | \psi_{ij}^* A(\mathcal{B}) \sum_{i'} \sum_{j'} \psi_{i'j'} |i'\rangle |j'\rangle$$

$$= \sum_{i'} \sum_{i''} \sum_{j'} \sum_{j''} \psi_{i'j'} \psi_{i''j''}^* \langle i' | A(\mathcal{B}) |i''\rangle \underbrace{\langle j' | j'' \rangle}_{\delta_{jj'}}$$

$$= \sum_{i'} \sum_{i''} \underbrace{\left( \sum_{j'} \psi_{i'j'} \psi_{i''j'}^* \right)}_{\rho_{i'i''}} \underbrace{\langle i' | A(\mathcal{B}) |i''\rangle}_{A_{i'i''}}$$

$$\circ \quad \therefore \langle A(\beta) \rangle = \sum_{i=1}^N \sum_{i'=1}^M P_{ii'} A_{ii'} = \text{tr}_{\beta}(PA) \quad (7)$$

where

$$P_{ii'} = \sum_{j=1}^M \psi_{ij} \psi_{i'j}^* \quad (8)$$

$$A_{ii'} = \langle i' | A(\beta) | i \rangle \quad (9)$$

The reduced matrix  $P_{ii'}$ , which operates for a block in environment, is the central vehicle of DMRG.

— Problem

Optimal  $P_{ii'}$  of rank- $m$  ( $\ll N$ )?

○ — Solution

Singular value decomposition (SVD) of  $\psi_{ij}$

## ○ — Singular value decomposition

Let assume  $N \geq M$ . Otherwise, we can simply switch the role of  $N$  and  $M$ .

(Theorem) An  $N \times M$  matrix  $\Psi$  ( $N \geq M$ ) can be decomposed as

$$N \begin{bmatrix} M \\ \Psi \end{bmatrix} = N \begin{bmatrix} M \\ U \end{bmatrix} \cdot M \begin{bmatrix} d_1 & & \\ & \ddots & \\ & & d_M \end{bmatrix} \cdot M \begin{bmatrix} M \\ V^T \end{bmatrix} \quad (10)$$

$$\text{or} \quad \begin{matrix} NM \\ \Psi \end{matrix} = \begin{matrix} NM & MM & MM \\ U & D & V^T \end{matrix} \quad (11)$$

where

$$D = \text{diag}(d_1, \dots, d_M) \quad (12)$$

$U = [u_i^{(v)}]$  is an  $N \times M$  column orthogonal matrix

$$\sum_{i=1}^N u_i^{(v)} u_i^{(v')} = \sum_{i=1}^N U_{vi}^T U_{iv'} = \delta_{vv'} \quad (13)$$

or

$$U^T U = I_M \quad (14)$$

and  $V = [v_j^{(v)}]$  is an  $M \times M$  column orthogonal matrix

$$\sum_{j=1}^M v_j^{(v)} v_j^{(v')} = \sum_{j=1}^M V_{vj}^T V_{jv'} = \delta_{vv'} \quad (15)$$

or

$$V^T V = I_M \quad (\text{orthonormality}) \quad (16)$$

Since  $V$  is square, it is also row-orthogonal,

$$V V^T = I_M \quad (\text{completeness}) \quad (17)$$

or

$$\sum_{v=1}^M v_j^{(v)} v_{j'}^{(v)} = \delta_{jj'} \quad (18)$$

⊙ (completeness)

Any vector  $a \in \mathbb{R}^M$  can be expressed as a linear combination of  $M$  linearly-independent vectors,

$$a_i = \sum_{\nu=1}^M c_{\nu} v_i^{(\nu)}$$

$$\sum_{i=1}^M v_i^{(\mu)} \times (\text{above})$$

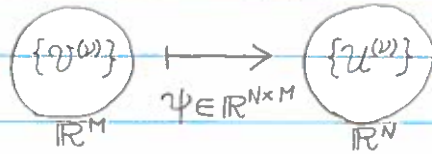
$$\sum_{i=1}^M v_i^{(\mu)} a_i = \sum_{\nu=1}^M c_{\nu} \underbrace{\sum_{i=1}^M v_i^{(\mu)} v_i^{(\nu)}}_{(V^T V)_{\mu\nu} = \delta_{\mu\nu}} = c_{\mu}$$

$$\begin{aligned} \therefore \forall a_i &= \sum_{\nu=1}^M v_i^{(\nu)} \underbrace{\sum_{i'=1}^M v_{i'}^{(\nu)} a_{i'}}_{c_{\nu}} \\ &= \sum_{i'=1}^M \underbrace{\left( \sum_{\nu=1}^M v_i^{(\nu)} v_{i'}^{(\nu)} \right)}_{\delta_{ii'}} a_{i'} \end{aligned}$$

$$\therefore (V V^T)_{ii'} = \sum_{\nu=1}^M v_i^{(\nu)} v_{i'}^{(\nu)} = \delta_{ii'} \quad //$$

(Interpretation of SVD)

In SVD,  $\psi \in \mathbb{R}^{N \times M}$  maps an orthonormal basis  $\{v^{(\mu)} \in \mathbb{R}^M\}$  onto another orthonormal basis  $\{u^{(\mu)} \in \mathbb{R}^N\}$ :



$$\psi v^{(\nu)} = u^{(\nu)} d_\nu \quad (\text{eigen-like relation}) \quad (19)$$

where

$$\begin{cases} v^{(\mu)T} v^{(\nu)} = \delta_{\mu\nu} \\ u^{(\mu)T} u^{(\nu)} = \delta_{\mu\nu} \end{cases} \quad (\text{orthonormality}) \quad (20)$$

There can be at maximum  $M$  such relations, since there are only  $M$  linearly-independent vectors in  $\mathbb{R}^M$ .

○ — Low-rank approximation to  $\rho_{ii'}$

(Theorem) Let  $\psi = UDV^T$  be the SVD of  $\psi$  with the diagonal elements in descending order  $d_1 \geq d_2 \geq \dots \geq d_M$ . Let

$$\psi^{(m)} = \sum_{\nu=1}^m \underset{N \times 1}{u^{(\nu)}} d_{\nu} \underset{1 \times M}{v^{(\nu)T}} \quad (21)$$

be the rank- $m$  truncation of the SVD. Then,

$$\min_{\text{rank}(A)=m} \|\psi - A\|_2 = \|\psi - \psi^{(m)}\|_2 = d_{m+1} \quad (22)$$

where the matrix 2-norm  $\|A\|_2 = \max_{\substack{\|x\|_2=1 \\ x \in \mathbb{R}^M}} \|Ax\|$  is defined in terms of vector 2-norms.

(White used the Frobenius norm,  $\|A\|_F = \sqrt{\sum_{ij} |a_{ij}|^2}$ .)

Therefore,  $\psi^{(m)}$  is the optimal rank- $m$  approximation to  $\psi$ .

○ (Reduced density matrix)

Substituting the rank- $m$  approximation Eq.(21) in Eq.(8),

$$\begin{aligned} \rho &= \psi \psi^T \\ &\simeq \sum_{\nu=1}^m u^{(\nu)} d_{\nu} v^{(\nu)T} \sum_{\nu'=1}^m v^{(\nu')} d_{\nu'} u^{(\nu')T} \\ &= \sum_{\nu} \sum_{\nu'} \underbrace{u^{(\nu)} d_{\nu} (v^{(\nu)} I v^{(\nu')}) d_{\nu'} u^{(\nu')T}}_{\delta_{\nu\nu'} (\odot \text{ orthogonality})} \\ &= \sum_{\nu=1}^m u^{(\nu)} d_{\nu}^2 u^{(\nu)T} \end{aligned}$$

$$\therefore \rho \simeq \sum_{\nu=1}^m u^{(\nu)} w_{\nu} u^{(\nu)T} \quad (23)$$

or

$$\rho_{ii'} \simeq \sum_{\nu=1}^m u_i^{(\nu)} w_{\nu} u_{i'}^{(\nu)} \quad (24)$$

where

$$w_{\nu} = d_{\nu}^2 \quad (25)$$

(Sum rule)

$$\text{Tr} \rho = \text{Tr} \psi \psi^T = \text{Tr} \psi^T \psi \quad (26)$$

Note that

$$\begin{aligned} \langle \psi | \psi \rangle &= \sum_{ij} \langle j | \langle i | \psi_{ij}^* \sum_{i'j'} \psi_{i'j'} | i' \rangle | j \rangle \\ &= \sum_{ij} \sum_{i'j'} \underbrace{\langle j | j \rangle}_{\delta_{ij'}} \underbrace{\langle i | i \rangle}_{\delta_{ii'}} \psi_{ij}^* \psi_{i'j'} \\ &= \sum_{ij} \psi_{ij}^* \psi_{ij} = \text{Tr} \psi^T \psi = 1 \end{aligned} \quad (27)$$

Therefore,

$$\text{Tr} \rho = \text{Tr} \psi \psi^T = 1 \quad (28)$$

Note that

$$\begin{aligned} \text{Tr} \rho &= \text{Tr} \psi \psi^T \\ &= \text{Tr} U \underbrace{D V^T V D}_{I_M} U^T \\ &= \text{Tr} U D^2 U^T \\ &= \text{Tr} \underbrace{U^T U}_{I_{MM}} D^2 \\ &= \text{Tr} D^2 \end{aligned} \quad (29)$$

Comparing Eqs. (28) and (29),

$$\text{Tr} \rho = \sum_{\nu=1}^M d_{\nu}^2 = \sum_{\nu=1}^M \omega_{\nu} = 1 \quad (30)$$

(Summary)

The rank- $m$  truncation of the SVD of the ground-state wavefunction,

$$\psi^{(m)} = \sum_{\nu=1}^m u^{(\nu)} d_{\nu} v^{(\nu)T} \quad (31)$$

or

$$\psi_{ij}^{(m)} = \sum_{\nu=1}^m u_i^{(\nu)} d_{\nu} v_j^{(\nu)} \quad (32)$$

produces the rank- $m$  approximation to the reduced density matrix,

$$\rho^{(m)} = \sum_{\nu=1}^m u^{(\nu)} w_{\nu} u^{(\nu)T} \quad (33)$$

or

$$\rho_{ii'}^{(m)} = \sum_{\nu=1}^m u_i^{(\nu)} w_{\nu} u_{i'}^{(\nu)} \quad (34)$$

where

$$w_{\nu} = d_{\nu}^2 \quad (35)$$

The rank- $m$  approximation  $\rho^{(m)}$  is optimal in the least square sense with the error estimate,

$$\text{Tr}(\rho - \rho^{(m)}) = 1 - \sum_{\nu=1}^m w_{\nu} \quad (36)$$



(Optimal reduced density matrix algorithm)

Given the superblock ground state  $\psi_{ij} \in \mathbb{R}^{N \times M}$

1. Obtain the  $N \times N$  reduced density matrix  $\rho_{ii'} \in \mathbb{R}^{N \times M}$  by integrating out the environment variables,

$$\rho_{ii'} = \sum_{j=1}^M \psi_{ij} \psi_{i'j}^* \quad (37)$$

2. Perform the SVD of  $\rho_{ii'}$

$$\rho_{ii'} = \sum_{\nu=1}^N u_i^{(\nu)} w_{\nu} u_{i'}^{(\nu)} \quad (38)$$

$$\text{or} \quad \rho = \sum_{\nu=1}^N u^{(\nu)} w_{\nu} u^{(\nu)T} \quad (39)$$

where  $w_1 \geq w_2 \geq \dots \geq w_N \geq 0$ .

3. Retain only the first  $m$  terms in the expansion

$$\rho_{ii'}^{(m)} = \sum_{\nu=1}^m u_i^{(\nu)} w_{\nu} u_{i'}^{(\nu)} \quad (40)$$

$$\text{or} \quad \rho^{(m)} = \sum_{\nu=1}^m u^{(\nu)} w_{\nu} u^{(\nu)T} \quad (41)$$

# Density Matrix Renormalization Group Revisited (2)

## Greedy / Growth Algorithm

6/17/03

[S.R. White, Phys. Rev. B 48, 10345 ('93); S. Daul, I. Ciofini, C. Daul, S.R. White, Intl. J. Quantum Chem. 79, 331 ('00)]

### — Objective

Obtain the ground state without diagonalizing a large matrix.

### — Approach: Greedy / growth algorithm

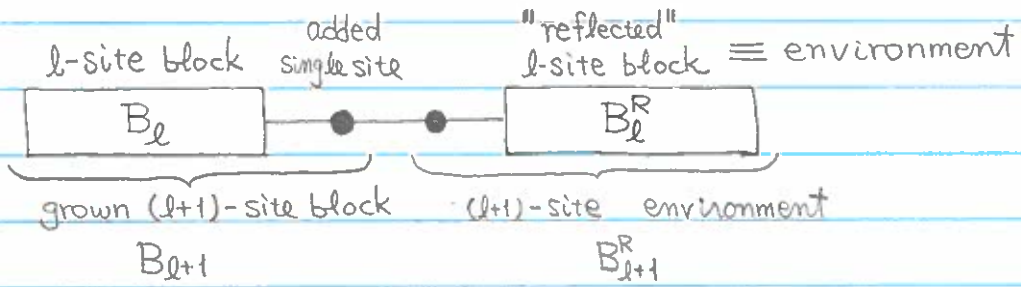
Grow an incrementally larger block (cf. Dijkstra's shortest-path algorithm), for which the optimal constant low-rank (rank- $m$ ) approximation is known (greedy).



By successively increasing the block size, while keeping the constant ( $m$ ) rank for its density matrix:

- i) Throw out higher excited states, which do not alter the ground state;
- ii) Add low-energy perturbation to refine the constant (rank- $m$ ) size renormalized basis.  
(dressed)

- Block/environment construction for 1D system.



- i) Reflected (right interchanged with left) block acts as environment.
- ii) Growth achieved by adding a single site.

- Infinite system algorithm

Starting from a superblock-environment consisting 4 single sites, grow the superblock each time by adding a single site until the superblock-environment contains  $L$  sites after  $(\frac{L}{2}-1)$  steps.

☺  $2(s-1) + 4 = L$   
 except for the 1st step, added per step  
 initially 4-site superblock+environment

$\therefore s-1 + 2 = \frac{L}{2}$   
 $\therefore s = \frac{L}{2} - 1$  //

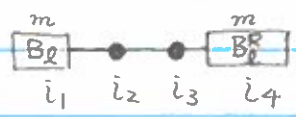
Step	Sites
1	4
2	6
3	8
...	
$\frac{L}{2} - 1$	$L$

Diagram for step  $\frac{L}{2} - 1$  shows a long horizontal bar representing the superblock-environment of length  $L$ , with a double-headed arrow below it indicating the total length.

1. Make four initial blocks, each consisting of a single site. Set up the block Hamiltonian.



2. Form the superblock Hamiltonian by adding cross terms.



3. Obtain the superblock ground state by the Davidson method,  $\Psi(i_1, i_2, i_3, i_4)$  (Compute physical expectation values here.)

4. Form the reduced density matrix for the 1-2 block by integrating out the environment variables (3-4),

$$\rho(i_1, i_2; i'_1, i'_2) = \sum_{i_3, i_4} \Psi(i_1, i_2, i_3, i_4) \Psi^*(i'_1, i'_2, i_3, i_4)$$

5. Diagonalize  $\rho$  to find eigenvalues  $w_\nu$  ( $w_1 \geq w_2 \geq \dots$ ) and eigenvectors  $U_{i_1 i_2}^{(\nu)}$ ; discard all but  $m$  largest eigenvalues.
6. Form matrix representation of operators for the two-block (1-2) system, and form a new block 1 by changing basis to  $U^{(\nu)}$ :  $H'_1 = O H_{12} O^T$ , where  $O_{\nu; i_1, i_2} = U_{i_1 i_2}^{(\nu)}$  (e.g. 1 site can add  $m_1 m_2$  degrees of freedom, and  $m_1 m_2 = 4m^2 \rightarrow m^2$  reduction).
7. Replace old block 1 with new block 1; replace old block 4 with the reflection of new block 1.
8. Go to step 2

○ - Lesson.

Block-enlarging perturbation (cf. Cuppen divide-and-conquer secular equation) should involve only renormalized/dressed (low-rank) operators — e.g. through Wannier orbitals without seeing real-space meshes/linear combination of Wannier functions.

↓

Local relaxation  $\longleftrightarrow$  reduced subspace diagonalization cycle?

○

○